## NOMENCLATURE

$A, B=$ van Laar constants
$P=$ total pressure of mixture
$P_{i}^{\circ}=$ vapor pressure of pure component $i$
$S S=$ sum of squares
$x_{1}=$ mole fraction of component $i$ in the liquid phase
$\gamma=$ activity coefficient of component $i$ in the liquid phase

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# Thermodynamic Functions of Methane 

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#### Abstract

Standard-state thermodynamic functions for gaseous methane are reported for the temperature range $60^{\circ} \mathrm{K}$. to $5000^{\circ} \mathrm{K}$., taking into consideration low-temperature quantum effects, anharmonicity, centrifugal distortion, and vibration-rotation interaction.


ALTHOUGH METHANE is a thermodynamically important molecule, no tables of its thermodynamic functions appear to have been published in which the rigid-rotator, harmonic oscillator ( RRHO ) approximation was abandoned. The best available values (11) are based on an RRHO calculation performed in 1944 (10). Spectroscopic information on methane has increased to the point where now the higher-order corrections can be applied. This paper reports the standard state thermodynamic functions for gaseous methane from $60^{\circ} \mathrm{K}$. to $5000^{\circ} \mathrm{K}$., taking into consideration low-temperature quantum effects, anharmonicity (16), centrifugal distortion (8), and vi-bration-rotation interaction (9).

The functions were calculated for an ideal gas at 1 atm. pressure. The RRHO values were obtained using the Viney (15) rotational partition function. Correction terms were included for anharmonicity through the second order (i.e., terms varying as $T^{2}$ at high temperatures). The contribution of nuclear spin was not included. Separate calculations were made for $\mathrm{C}^{12} \mathrm{H}_{4}$ and $\mathrm{C}^{13} \mathrm{H}_{4}$ using $\mathrm{C}^{13} \mathrm{H}_{4}$ frequencies and anharmonicities calculated as previously described (7). The maximum difference in $C_{p}^{0} / R$ between these calculations and the values obtained using $\mathrm{C}^{12} \mathrm{H}_{4}$ frequencies and anharmonicities and a mean molecular weight (based on $1.11^{\circ} \mathrm{c} \mathrm{C}^{13} \mathrm{H}_{i}$ ) was only 0.0002 . Since this is less than the probable error introduced by uncertainties in the molecular constants, the values reported here are based on the $\mathrm{C}^{12} \mathrm{H}_{4}$ constants and the mean molecular weight. Physical-scale constants (3) were used with $h c / k=1.43880 \mathrm{~cm}$. deg. and a translational entropy constant of -1.16552 . Calculations were carried out on an IBM 7090 computer.

The molecular constants selected for this work are listed in Table I.

Because of the extreme complexity of the methane spectrum, due in large part to the high degeneracy of most bands and the resultant opportunity for various types of interaction, the selection of a suitable set of constants is difficult. This will probably be the situation for some time to come, since the finer points of the spectra of tetrahedral molecules are only now beginning to be studied in detail $(4,5)$ and the effect of higher-order perturbations of such molecules on their thermodynamic functions

| Table I. Molecular Constants of Methane |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: |
| Mass $=16.04757 \mathrm{amu} .($ physical scale ) (14) |  |  |  |  |
| Rotational Constants (13) |  |  |  |  |
| $B_{6}=5.2406 \mathrm{~cm} .^{-1}$ |  |  |  |  |
| $D_{5}=1.23 \times 10^{-4} \mathrm{~cm} .^{-1}$ |  |  |  |  |
| $\sigma=$ |  |  |  |  |
| Vibrational Constants ( $\mathrm{cm} .^{-1}$ ) |  |  |  |  |
|  | $i=1$ | $i=2$ | $i=3$ | $i=4$ |
| $\nu:(2,13)$ | 2916.7 | 1533.6 | 3018.9 | 1305.9 |
| $X_{14}$ (7) | -64.6 | 0 | -65.0 | 0 |
| $X_{2 i}$ |  | 0 | -15.0 | -11.2 |
| $X_{3 i}$ |  |  | -17.5 | -12.0 |
| $\chi_{*}$ |  |  |  | -6.0 |
| $\alpha_{2}$ | -0.025 | -0.11 | 0.037 | 0.045 |

Table II. Thermodynamic Functions of Methane

| T, ${ }^{\circ} \mathrm{K}$. | $C_{p}^{o} / R$ | $\begin{gathered} \left(H^{\circ}-H_{o}^{o}\right) / \\ R T \end{gathered}$ | $\begin{gathered} -\left(F^{\circ}-H_{o}^{\circ}\right) / \\ R T \end{gathered}$ | $S^{\circ} / R$ |
| :---: | :---: | :---: | :---: | :---: |
| 60 | 4.0014 | 3.9693 | 11.9644 | 15.9337 |
| 80 | 4.0019 | 3.9774 | 13.1075 | 17.0849 |
| 90.67 | 4.0021 | 3.9803 | 13.6056 | 17.5859 |
| 100 | 4.0024 | 3.9823 | 13.9956 | 17.9779 |
| 111.67 | 4.0027 | 3.9844 | 14.4353 | 18.4197 |
| 120 | 4,0029 | 3.9857 | 14.7220 | 18.7077 |
| 140 | 4.0042 | 3.9882 | 15.3366 | 19.3248 |
| 160 | 4.0075 | 3.9904 | 15.8693 | 19.8597 |
| 180 | 4.0153 | 3.9927 | 16.3394 | 20.3321 |
| 200 | 4.0309 | 3.9956 | 16.7602 | 20.7559 |
| 220 | 4.0571 | 4.0000 | 17.1412 | 21.1412 |
| 240 | 4.0966 | 4.0063 | 17.4895 | 21.4958 |
| 260 | 4.1506 | 4.0152 | 17.8105 | 21.8257 |
| 273.15 | 4.1943 | 4.0227 | 18.0088 | 22.0316 |
| 280 | 4.2196 | 4.0272 | 18.1085 | 22.1357 |
| 298.15 | 4.2947 | 4.0412 | 18.3619 | 22.4030 |
| 300 | 4.3030 | 4.0428 | 18.3869 | 22.4296 |
| 320 | 4.3996 | 4.0620 | 18.6484 | 22.7104 |
| 340 | 4.5077 | 4.0850 | 18.8953 | 22.9803 |
| 360 | 4.6257 | 4.1117 | 19.1295 | 23.2412 |
| 380 | 4.7516 | 4.1420 | 19.3526 | 23.4947 |
| 400 | 4.8839 | 4.1758 | 19.5659 | 23.7417 |
| 450 | 5.2328 | 4.2737 | 20.0632 | 24.3369 |
| 500 | 5.5941 | 4.3876 | 20.5193 | 24.9069 |
| 550 | 5.9565 | 4.5138 | 20.9433 | 25.4571 |
| 600 | 6.3136 | 4.6489 | 21.3418 | 25.9907 |
| 650 | 6.6620 | 4.7904 | 21.7195 | 26.5099 |
| 700 | 6.9996 | 4.9363 | 22.0798 | 27.0160 |
| 750 | 7.3253 | 5.0847 | 22.4254 | 27.5101 |
| 800 | 7.6382 | 5.2346 | 22.7584 | 27.9930 |
| 850 | 7.9380 | 5.3849 | 23.0802 | 28.4651 |
| 900 | 8.2243 | 5.5347 | 23.3922 | 28.9270 |
| 950 | 8.4972 | 5.6835 | 23.6955 | 29.3790 |
| 1000 | 8.7566 | 5.8308 | 23.9908 | 29.8215 |
| 1050 | 9.0029 | 5.9760 | 24.2788 | 30.2548 |
| 1100 | 9.2363 | 6.1189 | 24.5601 | 30.6790 |
| 1150 | 9.4573 | 6.2593 | 24.8352 | 31.0945 |
| 1200 | 9.6663 | 6.3970 | 25.1045 | 31.5015 |
| 1300 | 10.0508 | 6.6635 | 25.6271 | 32.2906 |
| 1400 | 10.3944 | 6.9180 | 26.1303 | 33.0483 |
| 1500 | 10.7015 | 7.1601 | 26.6160 | 33.7761 |
| 1600 | 10.9767 | 7.3902 | 27.0855 | 34.4757 |
| 1700 | 11.2239 | 7.6086 | 27.5401 | 35.1487 |
| 1800 | 11.4469 | 7.8157 | 27.9809 | 35.7967 |
| 1900 | 11.6486 | 8.0122 | 28.4088 | 36.4211 |
| 2000 | 11.8320 | 8.1987 | 28.8246 | 37.0233 |
| 2100 | 11.9995 | 8.3758 | 29.2289 | 37.6047 |
| 2200 | 12.1531 | 8.5440 | 29.6225 | 38.1665 |
| 2300 | 12.2945 | 8.7040 | 30.0058 | 38.7098 |
| 2400 | 12.4255 | 8.8564 | 30.3795 | 39.2359 |
| 2500 | 12.5471 | 9.0016 | 30.7440 | 39.7456 |
| 2600 | 12.6607 | 9.1401 | 31.0998 | 40.2399 |
| 2700 | 12.7672 | 9.2725 | 31.4472 | 40.7197 |
| 2800 | 12.8675 | 9.3991 | 31.7867 | 41.1859 |
| 2900 | 12.9622 | 9.5203 | 32.1187 | 41.6390 |
| 3000 | 13.0521 | 9.6365 | 32.4434 | 42.0800 |
| 3200 | 13.2194 | 9.8553 | 33.0724 | 42.9277 |
| 3400 | 13.3732 | 10.0577 | 33.6761 | 43.7337 |
| 3600 | 13.5162 | 10.2458 | 34.2563 | 44.5021 |
| 3800 | 13.6507 | 10.4214 | 34.8150 | 45.2364 |
| 4000 | 13.7784 | 10.5860 | 35.3538 | 45.9398 |
| 4200 | 13.9007 | 10.7408 | 35.8741 | 46.6149 |
| 4400 | 14.0187 | 10.8870 | 36.3772 | 47.2642 |
| 4600 | 14.1333 | 11.0256 | 36.8642 | 47.8898 |
| 4800 | 14.2453 | 11.1573 | 37.3362 | 48.4935 |
| 5000 | 14.3552 | 11.2828 | 37.7943 | 49.0771 |

Table III. Contributions of Individual Terms to $C_{p}^{o} / R$

|  | $298.15^{\circ} \mathrm{K}$. | $600^{\circ} \mathrm{K}$. | $1000^{\circ} \mathrm{K}$. | $5000^{\circ} \mathrm{K}$. |
| :--- | :---: | :---: | :---: | ---: |
| RRHO |  |  |  |  |
| Ampproximation | 4.2866 | 6.2819 | 8.6352 | 12.6944 |
| Anharmonicity <br> Centrifugal | 0.0002 | 0.0197 | 0.1026 | 1.5114 |
| Distortion <br> Rot-Vib | 0.0071 | 0.0145 | 0.0248 | 0.1661 |
| Interaction | 0.0008 | 0.0025 | -0.0060 | -0.0167 |
| Total $C_{p}^{o} / R$ | 4.2947 | 6.3136 | 8.7566 | 14.3552 |

remains unknown. The treatment used here, however, should give a close approximation to the true thermodynamic functions and thus be a considerable improvement over the RRHO approximation.

The rotational constants obtained by Thomas and Welsh (13) from the fundamental $\nu_{3}$ were selected as being the best available. The rotation-vibration interaction constants, defined by $B_{v}=B_{0}-\sum \alpha_{i} v_{i}$, present particular problems, among which is the probability that this definition is something of an oversimplification for spherical top molecules; we shall assume that the relationship is valid at least for the lower vibrational states. Hecht (6) has estimated that $\alpha_{3}=0.042 \mathrm{~cm} .^{-1}$, but the overtone $2 \nu_{3}$ suggests a somewhat lower value $(4,12)$. An estimate of $\alpha_{2}$ is available from the work of Thomas and Welsh (13), and Fox's (4) recent analysis of $2 \nu_{4}$ gives $\alpha_{4}=0.045 \mathrm{~cm} .^{-1}$. No direct information is available on $\alpha_{1}$, but a study (1) of the combination band $\nu_{1}+\nu_{4}$ suggests that $\alpha_{1}+\alpha_{4} \sim 0.02$, giving a value for $\alpha_{1}$ which is consistent with the shape of the $Q$ branch of $\nu_{1}$ as observed in the Raman spectrum (13).

The calculated values of the thermodynamic functions obtained from these data are listed in Table II. Table III indicates the values of the individual correction terms to $C_{p}^{\circ} / R$ at several temperatures; the corrections to the other functions are somewhat smaller. The effect of vibrationrotation interaction is greater than the figures of Table III suggest; the apparently small contributions listed there result from fortuitous near-cancellation of the terms due to $\alpha_{2}$ and $\alpha_{4}$, which are the major contributors. Actually the uncertainty in the $\alpha$ 's is the greatest source of uncertainty in these calculations; a change in $\alpha_{4}$ (the most sensitive constant) of $\pm 0.01 \mathrm{~cm} .^{-1}$, for example, will result in a change of $\pm 0.0090$ in $C_{p}^{o} / R$ at $1000^{\circ} \mathrm{K}$.

In a recently published paper, J. Herranz and B.P. Stoicheff, J. Mol. Spectr. 10, 448 (1963), have derived slightly different molecular constants from the $\nu_{3}$ Raman band: $\nu_{3}=3019, \mathrm{~B}_{o}=5.2412, \mathrm{D}_{o}=1.12 \times 10^{-4}, \alpha_{3}=0.0416 \mathrm{~cm} .^{-1}$. We have recalculated the thermodynamic functions using these values in place of the corresponding constants in Table I, and find that $C_{p}^{\circ} / R$ is raised $0.02 \%$ at $1000^{\circ} \mathrm{K}$. and $0.08 \%$ at $5000^{\circ} \mathrm{K}$.; the effect on the other functions is even less.

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